

An Asymptotic-Preserving Particle-In-Cell method for the Vlasov-Poisson system near quasineutrality.

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Saint-Malo the 27th january 2011

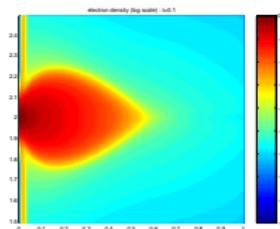
Joint works with

- P. Degond, F. Deluzet, L. Navoret, A.B. Sun
- C. Chainais-Hillairet

General topic : multiscale problems

- **Mathematic models, M_λ , depending on a parameter λ**

- λ very small in a part of the domain
 - λ is an order 1 parameter elsewhere



- **Two-scales :** microscopic scale (small values of λ)
and

macroscopic scale (large values of λ)

- **Difficulties :** The classical explicit schemes are stable and consistent

iff λ is resolved by the mesh \Rightarrow **huge cost**

A possible solution : use the asymptotic model where $\lambda \ll 1$

$$M_0 = \lim_{\lambda \rightarrow 0} M_\lambda \quad \Rightarrow \quad \text{Mesh independent of } \lambda$$

► Problems

- Boundary conditions **well prepared** to the asymptotic regime
- You must use M_λ where $\lambda = O(1)$ \Rightarrow **two different models**
 - Reconnection of M_λ and M_0 .
 - Location of the interface.
 - Moving interface : difficult numerical problem in 2D or 3D.

Another solution : use an asymptotic preserving scheme

- ➡ Use the microscopic model M_λ everywhere.
- ➡ Discretize it with a scheme which preserves the asymptotic limit $\lambda \rightarrow 0$
 - ➡ The mesh is independent of λ : **Asymptotic stability**.
 - ➡ You recover an approximate solution of M_0 when $\lambda \ll 1$:
Asymptotic consistency.

Scheme asymptotically stable and consistent

⇒ **Asymptotic preserving scheme (AP)**
([S.Jin, 1999] kinetic → Hydro)

Particular case of quasineutral limits

- ➡ In gaz containing charged particles

$$\lambda = \frac{\text{Debye length}}{\text{size of the domain}}$$

- ➡ Debye length is the scale of electric interactions in the gaz
- ➡ Quasineutral limits : $\lambda \rightarrow 0$
- ➡ Euler-Poisson :
- ➡ Vlasov-Poisson :
- ➡ Drift-diffusion Poisson :

► Quasineutral limits

► Euler-Poisson :

- S. Fabre (1992),
- Ph. Colella, M.R. Dorr, D.D. Wake (1999),
- H.H. Choe, N.S. Yoon, S.S. Kim, D.I. Choi (2001),
- P. Crispel, P. Degond, M.H. Vignal (2007),
- P. Degond, J.G. Liu, M.H. Vignal (2008),
- M.H. Vignal (2010).

► Vlasov-Poisson :

► Drift-diffusion Poisson :

➡ Quasineutral limits

➡ Euler-Poisson :

➡ Vlasov-Poisson :

- R.J. Mason (1981,1983),
- B.I. Cohen, A.B. Langdon, A. Friedman (1982,1983),
- P. Degond, F. Deluzet, L. Navoret (2006),
- R. Belaouar, N. Crouseilles, P. Degond, E. Sonnendrücker (2009),
- P. Degond, F. Deluzet, L. Navoret, A.B. Sun, M.H. Vignal (2010).

➡ Drift-diffusion Poisson :

➡ Quasineutral limits

➡ Euler-Poisson :

➡ Vlasov-Poisson :

➡ Drift-diffusion Poisson :

- G. Lapenta, F. Iinoya, J.U. Brackbill (1995),
- P.L.G. Ventzak, T.J. Sommerer, R.J. Hoekstra, M.J. Kushner, (1993),
- P.L.G. Ventzak, R.J. Hoekstra, M.J. Kushner, (1994),
- C. Chainais-Hillairet, M.H. Vignal.

Outline

- 1 General topic : multiscale problems
- 2 An AP scheme for the Vlasov-Poisson model in the quasineutral limit
 - The Vlasov-Poisson model and its quasineutral limit
 - The classical explicit scheme and our AP scheme
 - Numerical results
- 3 Generalization to the Drift-Diffusion Poisson system
 - The Drift-Diffusion Poisson system and its quasineutral limit
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 - Numerical results
- 4 Works in progress

The Vlasov-Poisson model

- Joint work with P. Degond, F. Deluzet, L. Navoret, A.B. Sun,
- One species model for clarity

$$(VP)_\lambda \begin{cases} \partial_t f + v \cdot \nabla_x f + \nabla_x \phi \cdot \nabla_v f = 0, \\ -\lambda^2 \Delta \phi = n_0 - n, \quad n = \int f dv. \end{cases}$$

- $f(x, v, t)$ = distribution function, $\phi(x, t)$ = electric potential,
 n_0 = constant ion density, $n(x, t)$ = electron density,

$$\lambda = \frac{\text{Debye length}}{\text{size of the domain}} = \text{rescaled Debye length}$$

What is the quasineutral limit $\lambda \rightarrow 0$

- ➡ The quasineutral Vlasov system

$$(QNV) \left\{ \begin{array}{l} \partial_t f + v \cdot \nabla_x f + \nabla_x \phi \cdot \nabla_v f = 0, \\ n = \int f dv = n_0. \end{array} \right.$$

ϕ is the Lagrange multiplier of the constraint $n = n_0$

- ➡ Rigorous quasineutral limits : Brenier, Brenier & Grenier, Brenier & Corrias
- ➡ Explicit equation for ϕ ?

The quasineutral limit $\lambda \rightarrow 0$

➡ Taking the velocity moments of Vlasov and using $n = n_0$

$$\left\{ \begin{array}{l} \partial_t n + \nabla_x \cdot (n u) = 0, \\ \partial_t (n u) + \nabla_x S = n \nabla_x \phi, \\ n u = \int f v dv, \quad S = \int f v \otimes v dv. \end{array} \right. \Leftrightarrow \left\{ \begin{array}{l} \nabla_x \cdot (n u) = 0, \\ \partial_t (n u) + \nabla_x S = n \nabla_x \phi, \\ n u = \int f v dv, \quad S = \int f v \otimes v dv. \end{array} \right. \quad (1)$$

➡ Quasineutral elliptic equation

$$\nabla_x \cdot \left(\partial_t (n u) + \nabla_x S = n \nabla_x \phi \right) \quad \Rightarrow \quad \nabla_x \cdot (n u) = 0$$

$$-\nabla_x \cdot (n \nabla_x \phi) = -\nabla_x^2 : S$$

➡ Reciprocally

$$-\nabla_x \cdot (n \nabla_x \phi) = -\nabla_x^2 : S \quad \Rightarrow \quad \partial_t \nabla_x \cdot (n u) = 0 \quad \Rightarrow \quad \partial_{tt}^2 n = 0$$

The quasineutral limit $\lambda \rightarrow 0$

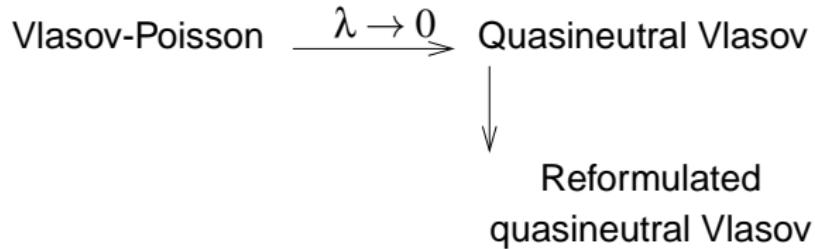
- ➡ Reformulated quasineutral Vlasov system

$$(RQNV) \left\{ \begin{array}{l} \partial_t f + v \cdot \nabla_x f + \nabla_x \phi \cdot \nabla_v f = 0, \\ -\nabla_x \cdot (n \nabla_x \phi) = -\nabla_x^2 : S, \\ (n = n_0)|_{t=0} \text{ and } \frac{d}{dt}(n = n_0)|_{t=0}, \\ n = \int f dv, \quad S = \int f v \otimes v dv. \end{array} \right.$$

- ➡ Equivalent to the quasineutral Vlasov model

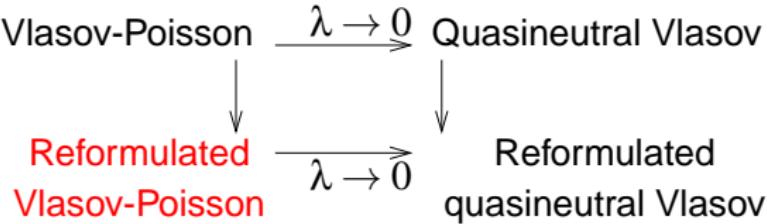
The reformulated Vlasov-Poisson system

➡ Is it possible to complete the diagram ?



The reformulated Vlasov-Poisson system

- It is possible to complete the diagram.



- Starting from Vlasov-Poisson and taking the velocity moments

$$\left\{ \begin{array}{l} \partial_t n + \nabla_x \cdot (n u) = 0, \\ \partial_t(n u) + \nabla_x S = n \nabla_x \phi. \end{array} \right. \quad \begin{array}{l} (2) \\ (3) \end{array}$$

$$\nabla_x \cdot (3) - \partial_t (2) \Rightarrow -\partial_{tt}^2 n - \nabla_x \cdot (n \nabla_x \phi) = -\nabla_x^2 : S$$

$$n = n_0 + \lambda^2 \Delta \phi \Rightarrow$$

The reformulated Poisson equation

$$\lambda^2 \partial_{tt}^2 (-\Delta \phi) - \nabla_x \cdot (n \nabla_x \phi) = -\nabla_x^2 : S$$

The reformulated Vlasov-Poisson system

$$(RVP) \left\{ \begin{array}{l} \partial_t f + v \cdot \nabla_x f + \nabla_x \phi \cdot \nabla_v f = 0, \\ \lambda^2 \partial_{tt}^2 (-\Delta \phi) - \nabla_x \cdot (n \nabla_x \phi) = -\nabla_x^2 : S, \\ (\lambda^2 \Delta \phi = n - n_0)|_{t=0} \text{ and } \frac{d}{dt}(\lambda^2 \Delta \phi = n - n_0)|_{t=0}. \end{array} \right. \quad (4)$$

- ➡ Equivalent to the Vlasov-Poisson model.
- ➡ Does not degenerate when $\lambda \rightarrow 0$.
- ➡ Reduces to (RQN) system when $\lambda = 0$. Consistency
with $\lambda \rightarrow 0$
- ➡ (4) is a harmonic oscillator Eq. on $-\Delta \phi$
- ➡ Oscillations in time at the period λ
 - ➡ Explicit scheme \Rightarrow conditional stab. $\Delta t \leq \lambda$
 - ➡ Implicit scheme \Rightarrow unconditional stability

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The classical scheme, semi-discretization in time

► If f^m approximates f at time t^m ($n^m = \int f^m dv$, $(nu)^m = \int f^m v dv$)

$$\begin{cases} \frac{f^{m+1} - f^m}{\Delta t} + v \cdot \nabla_x f^m + \nabla_x \phi^{m+1} \cdot \nabla_v f^m = 0, \\ -\lambda^2 \Delta \phi^{m+1} = n_0 - n^{m+1}. \end{cases}$$

► Uncoupled scheme :

► Integrating Vlasov eq. $\Rightarrow \frac{n^{m+1} - n^m}{\Delta t} + \nabla_x \cdot (nu)^m = 0$ gives n^{m+1} ,

► Poisson gives ϕ^{m+1} ,

► Vlasov gives f^{m+1} .

The discret reformulated Poisson equation

$$-\lambda^2 \frac{\Delta \phi^{m+1} - 2\Delta \phi^m + \Delta \phi^{m-1}}{\Delta t^2} - \nabla_x \cdot (n^{m-1} \nabla_x \phi^m) = \dots$$

► Stable and consistent iff $\Delta t \leq \lambda$.

The AP scheme, semi-discretization in time

► If f^m approximates f at time t^m ($n^m = \int f^m dv$, $(nu)^m = \int f^m v dv$)

$$\begin{cases} \frac{f^{m+1} - f^m}{\Delta t} + v \cdot \nabla_x f^{m+1} + \nabla_x \phi^{m+1} \cdot \nabla_v f^m = 0, \\ -\lambda^2 \Delta \phi^{m+1} = n_0 - \tilde{n}^{m+1}. \end{cases}$$

► Uncoupled scheme ?

► Integrating Vlasov eq. \Rightarrow

$$\frac{n^{m+1} - n^m}{\Delta t} + \nabla_x \cdot (nu)^{m+1} = 0,$$

$$\frac{(nu)^{m+1} - (nu)^m dv}{\Delta t} + \nabla_x \int f^{m+1} v \otimes v dv = n^m \nabla_x \phi^{m+1},$$

$$n^{m+1} = n^m - \Delta t \nabla_x \cdot (nu)^m + \Delta t^2 \left(\nabla_x^2 : \int f^{m+1} v \otimes v dv - \nabla_x \cdot (n^m \nabla_x \phi^{m+1}) \right)$$

The AP scheme, semi-discretization in time

► If f^m approximates f at time t^m ($n^m = \int f^m dv$, $(nu)^m = \int f^m v dv$)

$$\begin{cases} \frac{f^{m+1} - f^m}{\Delta t} + v \cdot \nabla_x f^{\textcolor{red}{m+1}} + \nabla_x \phi^{m+1} \cdot \nabla_v f^m = 0, \\ -\lambda^2 \Delta \phi^{m+1} = n_0 - \tilde{n}^{m+1}. \end{cases}$$

► Uncoupled scheme

► Replace n^{m+1} by \tilde{n}^{m+1} where

$$\tilde{n}^{m+1} = n^m - \Delta t \nabla_x \cdot (nu)^m + \Delta t^2 \left(\nabla_x^2 : \int f^{\textcolor{red}{m}} v \otimes v dv - \nabla_x \cdot \left(n^m \nabla_x \phi^{m+1} \right) \right)$$

► Unconditionally stable and consistent $\Delta t \leq \lambda$

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4 Works in progress

- One dimensional two-species plasma expansion test case

- Vlasov for ions $\partial_t f_i + v \partial_x f_i - \partial_x \phi \partial_v f_i = 0,$

- Vlasov for electrons $\partial_t f_e + v \partial_x f_e + \frac{\partial_x \phi}{\epsilon} \partial_v f_e = 0,$

- Poisson $-\lambda^2 \partial_{xx}^2 \phi = n_i - n_e, \quad n_{i,e} = \int f_{i,e} dv.$

$$\epsilon = \frac{\text{electron mass}}{\text{ion mass}}, \quad \text{Debye length} = \lambda \ll 1, \quad \text{Plasma period} = \sqrt{\epsilon} \lambda \ll 1.$$

- Initially, ions and electrons are Maxwellian

- Ion slab \rightarrow Electrons in a Boltzmannian equilibrium

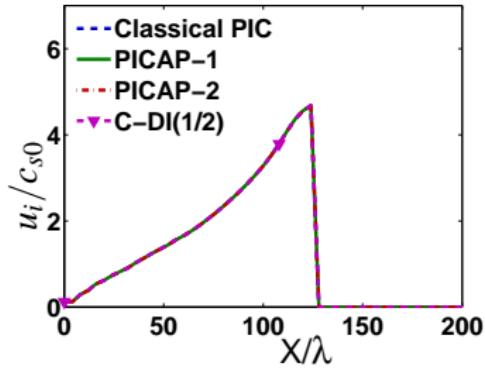
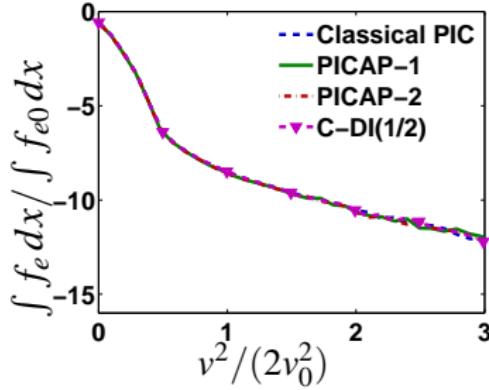
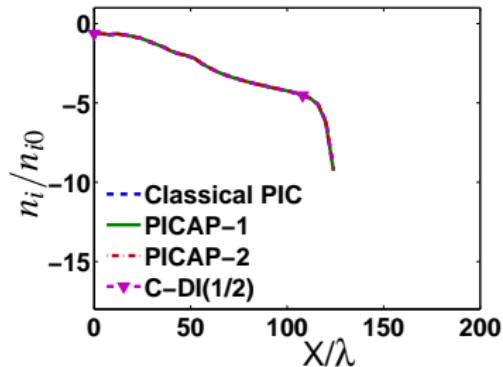
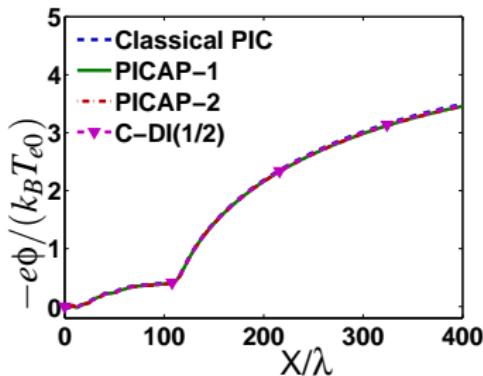
$$n_{i0} = \begin{cases} 1, & 0 \leq x \leq 20\lambda, \\ 0, & 20\lambda \leq x \leq 3.10^4. \end{cases}$$

$$\begin{cases} n_{e0} = \exp(\phi_0), \\ -\partial_{xx}^2 \phi_0 = n_{i0} - \exp(\phi_0). \end{cases}$$

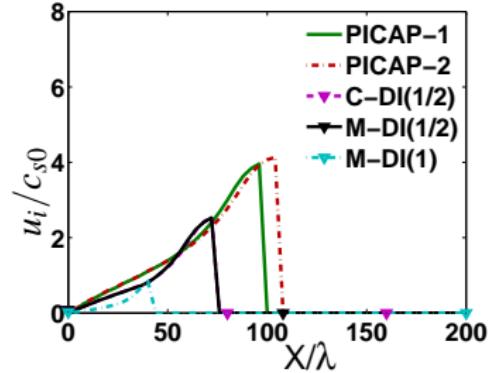
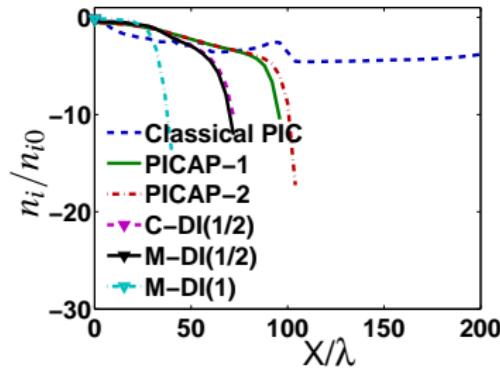
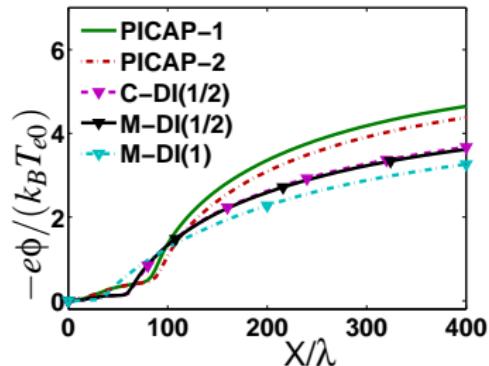
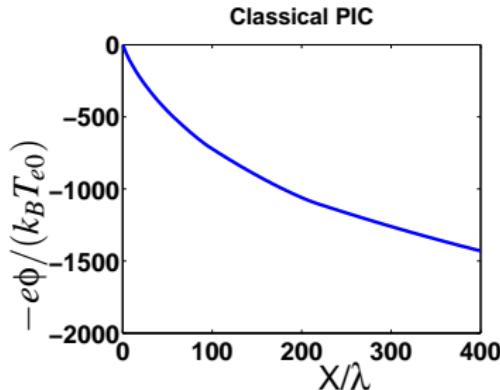
The test case

- ➡ Numerical parameters :
 - ➡ Domain : $x \in [0, 3.10^4 \lambda]$, with $\Delta x = 4\lambda \Rightarrow 7500$ cells.
 - ➡ Total number of numerical particles (ions+electrons) $\approx 5.10^6$.
 - ➡ Resolved case : $\Delta t = 0.05 \sqrt{\epsilon} \lambda$, Unresolved case : $\Delta t = 3 \sqrt{\epsilon} \lambda$
- ➡ Comparison between :
 - ➡ Explicit PIC scheme \Rightarrow conditional stability $\Delta t \leq \sqrt{\epsilon} \lambda$
 - ➡ Our AP-PIC scheme \Rightarrow unconditional stability $\Delta t \leq \sqrt{\epsilon} \lambda$
 - ➡ Direct-Implicit PIC schemes : Cohen, Friedman, Langdon $\Rightarrow \Delta t \leq \sqrt{\epsilon} \lambda$
(modified on the 2 first steps)

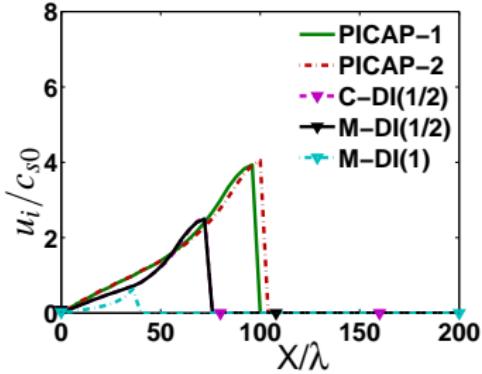
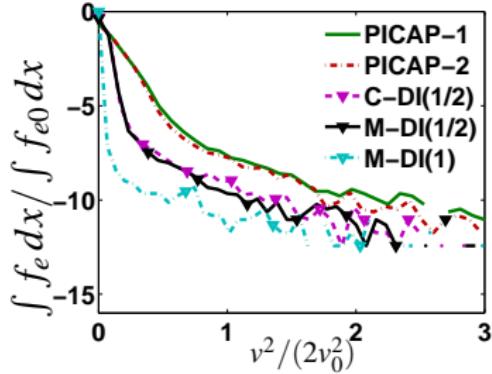
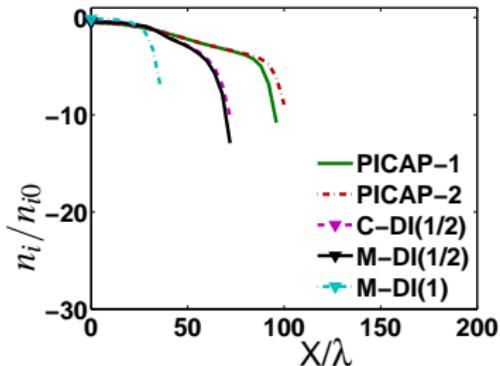
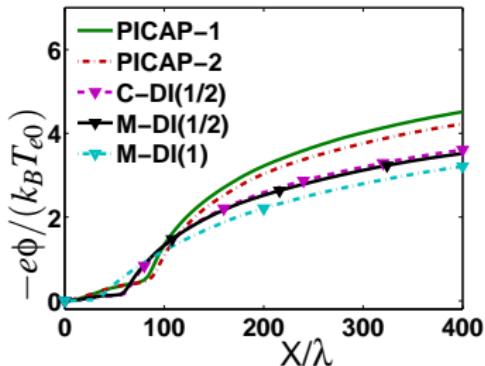
Resolved case : $\Delta t = 0.05 \sqrt{\varepsilon} \lambda$, $\Delta x = 4 \lambda$



Unresolved case : $\Delta t = 3 \sqrt{\varepsilon} \lambda$, $\Delta x = 4 \lambda$



Unresolved case, less particles /20



Gain of CPU time

➡ Ratios

$$\frac{\text{CPU Resolved case}}{\text{CPU Unresolved case}} = 48,$$

$$\frac{\text{CPU Resolved case}}{\text{CPU Unresolved case less particles}} = 960.$$

➡ About **1000 times faster** in one dimension.

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The scaled Drift-Diffusion Poisson system

- Joint work with C. Chainais-Hillairet (University Lille 1)

$$(DD)_\lambda \begin{cases} \partial_t N - \nabla \cdot (\nabla r(N) - N \nabla \phi) = 0, \\ \partial_t P - \nabla \cdot (\nabla r(P) + P \nabla \phi) = 0, \\ -\lambda^2 \Delta \phi = P - N + C, \end{cases} \quad r(s) = s^\gamma, \quad \gamma \geq 1.$$

- N = electron density, P = hole density,
 ϕ = electric potential, $C(x)$ = doping profile,

λ = rescaled Debye length

- Existence, sometimes uniqueness :

-T. Seidman, G. Troianello (1985), -H. Gajewski (1985),
-P.A. Markowitch, C.A. Ringhofer, C. Schmeiser (1990).

- Proof based on estimates in $L^\infty, L^2(0, T; H^1(\Omega))$ and Schauder fixed point

The formal quasineutral limit $\lambda \rightarrow 0$:

$$(DD)_0 \begin{cases} \partial_t N - \nabla \cdot (\nabla r(N) - N \nabla \phi) = 0, \\ \partial_t P - \nabla \cdot (\nabla r(P) + P \nabla \phi) = 0, \\ 0 = P - N + C. \end{cases}$$

- ➡ Subtracting mass eqs.

$$\partial_t (P - N) - \Delta(r(P) - r(N)) - \nabla \cdot ((P + N) \nabla \phi) = 0.$$

- ➡ Using the quasineutrality constraint \Rightarrow explicit eq. for the potential

$$-\nabla \cdot ((P + N) \nabla \phi) = \Delta(r(P) - r(N)).$$

- ➡ Starting from the drift-diffusion Poisson model

Same transformations \Rightarrow The reformulated Poisson equation

$$-\lambda^2 \partial_t \Delta \phi - \nabla \cdot ((P + N) \nabla \phi) = \Delta(r(P) - r(N)).$$

- ➡ Solutions oscillate at the period λ^2 .
- ➡ Explicit schemes must resolve this scale for stability $\Rightarrow \Delta t \leq \lambda^2$.

Rigorous quasineutral limit

➡ I. Gasser, C.D. Levermore, P.A. Markowich, C. Schmeiser (2001)

$$(DD)_\lambda \begin{cases} \partial_t N - \nabla \cdot (\nabla N - N \nabla \phi) = 0, \\ \partial_t P - \nabla \cdot (\nabla P + P \nabla \phi) = 0, \quad x \in \Omega \subset \mathbb{R}^d, t > 0. \\ -\lambda^2 \Delta \phi = P - N + C, \end{cases}$$

➡ Mixed quasineutral boundary conditions (Dirichlet and Neumann)

- On $\partial\Omega_D$: $N = N_b, \quad P = P_b, \quad \phi = \phi_b, \quad$ with $P_b - N_b + C = 0.$
- On $\partial\Omega_N$: $(\nabla N - N \nabla \phi) \cdot \mathbf{v} = 0, \quad (\nabla P + P \nabla \phi) \cdot \mathbf{v} = 0, \quad \nabla \phi \cdot \mathbf{v} = 0.$

➡ Quasineutral initial conditions

$$P_0 - N_0 + C = 0.$$

The rigorous quasineutral limit

Theorem : (Gasser-Levermore-Markowich-Schmeiser)

If - $C \in H^1(\Omega)$, - $P_0, N_0 \in L^1(\Omega)$ uniformaly in λ ,
- $N + P \geq C$, - Entropy (energy) $e(t = 0)$ is independant of λ .

Then $N^\lambda, P^\lambda \rightarrow N, P$ in $L^1([0, T] \times \Omega)$ and $\nabla \phi^\lambda \rightharpoonup \nabla \phi$ in $L^2(\Omega \times [0, T])$,

$$(DD)_0 \left\{ \begin{array}{l} \partial_t N - \nabla \cdot (\nabla N - N \nabla \phi) = 0, \\ \partial_t P - \nabla \cdot (\nabla P + P \nabla \phi) = 0, \\ 0 = P - N + C. \end{array} \right.$$

- On $\partial \Omega_D$: $N = N_b, P = P_b, \phi = \phi_b$.
- On $\partial \Omega_N$: $(\nabla N - N \nabla \phi) \cdot \mathbf{v} = 0, (\nabla P + P \nabla \phi) \cdot \mathbf{v} = 0$.

$$P_0 - N_0 + C = 0.$$

The rigorous quasineutral limit

► Choosing $P+N$, $P-N$ and ϕ for unknowns

$$(DD)_0 \Leftrightarrow \begin{cases} \partial_t(P+N) - \nabla \cdot (\nabla(P+N) - C\nabla\phi) = 0, \\ -\nabla \cdot (-\nabla C + (P+N)\nabla\phi) = 0, \\ P-N = -C. \end{cases}$$

- On $\partial\Omega_D$: $P+N = P_b + N_b$, $P-N = P_b - N_b = -C$, $\phi = \phi_b$.
- On $\partial\Omega_N$: $(\nabla(P+N) - C\nabla\phi) \cdot \mathbf{v} = 0$, $(-\nabla C + (P+N)\nabla\phi) \cdot \mathbf{v} = 0$.

► Proof based on :
 ► the entropy dissipation,
 ► Sobolev imbeddings.

► If $C = 0$, simpler, because estimates in L^∞ and $L^2(0, T; H^1(\Omega))$ are uniform in λ , (I. Gasser, 2001).

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The implicit scheme

- If N^n, P^n, ϕ^n are known approximations at time t^n

$$\begin{cases} N^{n+1} - N^n - \Delta t \nabla \cdot (\nabla N^{n+1} - N^{n+1} \nabla \phi^{n+1}) = 0, \\ P^{n+1} - P^n - \Delta t \nabla \cdot (\nabla P^{n+1} + P^{n+1} \nabla \phi^{n+1}) = 0, \\ -\lambda^2 \Delta \phi^{n+1} = P^{n+1} - N^{n+1} + C. \end{cases}$$

- Coupled scheme \Rightarrow non linear \Rightarrow Newton algorithm.
- Unconditionally stable scheme \Rightarrow asymptotic preserving scheme.

The implicit scheme, rigorous results

- ➡ Existence and uniqueness results
 - ➡ Semi-discretization in time : - estimates in $L^\infty, L^2(0, T; H^1(\Omega))$,
- Schauder fixed point.
 - ➡ Scheme in space and time : - L^∞ estimates, topological degree,
(C. Chainais-Hillairet, J.G. Liu, Y.J. Peng, 2003,2004).
- ➡ For $C = 0$, estimates in L^∞ and $L^2(]0, T[; H^1(\Omega))$ are uniform in λ .
- ➡ Discrete dissipation of the entropy (Marianne Chatard, to appear).
- ➡ For the semi-discretization in time : result similar to the convergence proof of Gasser-Levermore-Markowich-Schmeiser can be proved.

Our uncoupled asymptotic preserving scheme

➡ A scheme stable uniformly

➡ Implicit coupling terms \Rightarrow stability uniform in λ

$$\left\{ \begin{array}{l} N^{n+1} - N^n - \Delta t \nabla \cdot (\nabla N^n - N^n \nabla \phi^{n+1}) = 0, \\ P^{n+1} - P^n - \Delta t \nabla \cdot (\nabla P^n + P^n \nabla \phi^{n+1}) = 0, \\ -\lambda^2 \Delta \phi^{n+1} = P^{n+1} - N^{n+1} + C, \end{array} \right. \quad (1)$$

$$\left\{ \begin{array}{l} N^{n+1} - N^n - \Delta t \nabla \cdot (\nabla N^n - N^n \nabla \phi^{n+1}) = 0, \\ P^{n+1} - P^n - \Delta t \nabla \cdot (\nabla P^n + P^n \nabla \phi^{n+1}) = 0, \\ -\lambda^2 \Delta \phi^{n+1} = P^{n+1} - N^{n+1} + C, \end{array} \right. \quad (2)$$

$$\left\{ \begin{array}{l} N^{n+1} - N^n - \Delta t \nabla \cdot (\nabla N^n - N^n \nabla \phi^{n+1}) = 0, \\ P^{n+1} - P^n - \Delta t \nabla \cdot (\nabla P^n + P^n \nabla \phi^{n+1}) = 0, \\ -\lambda^2 \Delta \phi^{n+1} = P^{n+1} - N^{n+1} + C, \end{array} \right. \quad (3)$$

➡ Eliminating P^{n+1} and N^{n+1} , equation (3) yields

$$-\lambda^2 \Delta \phi^{n+1} - \Delta t \nabla \cdot ((N^n + P^n) \nabla \phi^{n+1}) = P^n - N^n + C + \Delta t \Delta (P^n - N^n).$$

➡ Discret reformulated Poisson equation

$$-\nabla \cdot ((\lambda^2 + \Delta t (N^n + P^n)) \nabla \phi^{n+1}) = P^n - N^n + C + \Delta t \Delta (P^n - N^n).$$

Our uncoupled asymptotic preserving scheme

- We replace in the implicit scheme the Poisson eq. by the reformulated Poisson eq.

$$\begin{cases} N^{n+1} - N^n - \Delta t \nabla \cdot (\nabla N^{n+1} - N^{n+1} \nabla \phi^{n+1}) = 0, \\ P^{n+1} - P^n - \Delta t \nabla \cdot (\nabla P^{n+1} + P^{n+1} \nabla \phi^{n+1}) = 0, \\ -\nabla \cdot ((\lambda^2 + \Delta t (N^n + P^n)) \nabla \phi^{n+1}) = P^n - N^n + C + \Delta t \Delta (P^n - N^n). \end{cases}$$

- Uncoupled scheme \Rightarrow linear.
- Asymptotic preserving scheme (verified in numerical results).

Our uncoupled asymptotic preserving scheme, results

➡ Existence and uniqueness of the solution : like for the implicit scheme

➡ Formally

⇒ $\lambda = 0$ in the scheme

$$\begin{cases} N^{n+1} - N^n - \Delta t \nabla \cdot (\nabla N^{n+1} - N^{n+1} \nabla \phi^{n+1}) = 0, \\ P^{n+1} - P^n - \Delta t \nabla \cdot (\nabla P^{n+1} + P^{n+1} \nabla \phi^{n+1}) = 0, \\ -\nabla \cdot ((\Delta t (N^n + P^n)) \nabla \phi^{n+1}) = P^n - N^n + C + \Delta t (P^n - N^n). \end{cases}$$

⇒ Asymptotic expansion in Δt

$\Delta t^0 \rightarrow$ The quasineutrality : $P^n - N^n + C = 0$,

$\Delta t^1 \rightarrow$ discretization of the quasineutral eq. for the potential ϕ :

$$-\nabla \cdot ((N^n + P^n) \nabla \phi^{n+1}) = \Delta (P^n - N^n) = -\Delta C.$$

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 - The Vlasov-Poisson model and its quasineutral limit
 - The classical explicit scheme and our AP scheme
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- 4 Works in progress

$\partial\Omega_D$

$$C = -1$$

$$C = 1$$

$\partial\Omega_D$

➡ $\lambda^2 = 10^{-3}$ or 10^{-10} , $\Delta t = 5 \cdot 10^{-4}$.

➡ Final time $T=0.5$

➡ Mesh in space : 896 triangles.

➡ Boundary conditions :

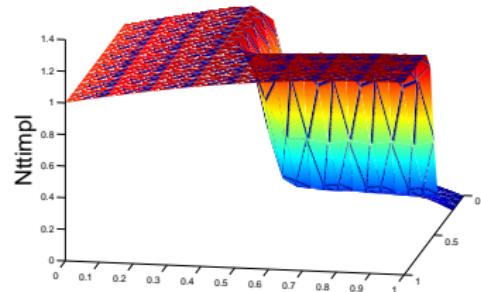
➡ On $\partial\Omega_D$: quasineutral Dirichlet conditions.

➡ Elsewhere : homogeneous Neumann conditions.

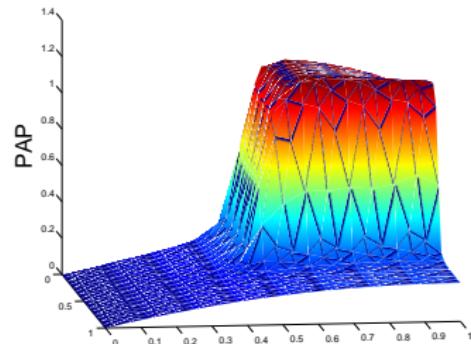
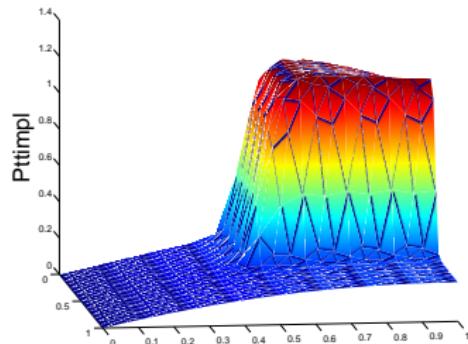
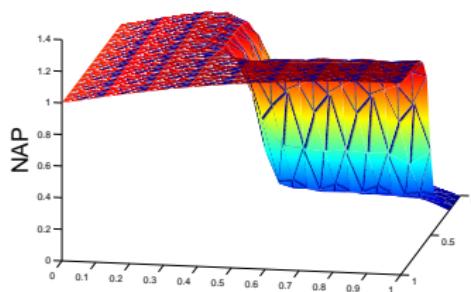
➡ Quasineutral initial conditions.

PN Diode, $\lambda^2 = 10^{-3}$, $\Delta t = 5 \cdot 10^{-4}$

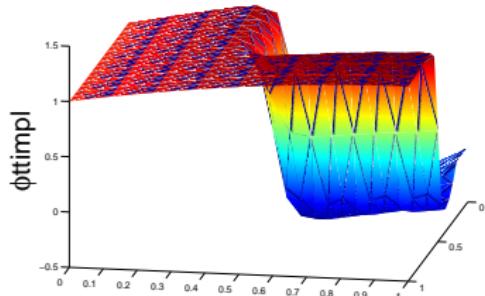
Implicit scheme



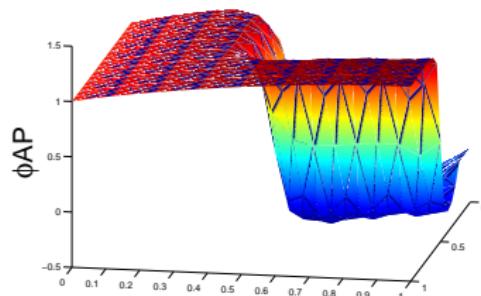
Our uncoupled AP scheme



Implicit scheme



Our uncoupled AP scheme



-CPU time of the implicit scheme : 86.7656s.

⇒ ratio ≈ 5.3 .

-CPU time of our uncoupled AP scheme : 16.5000s.

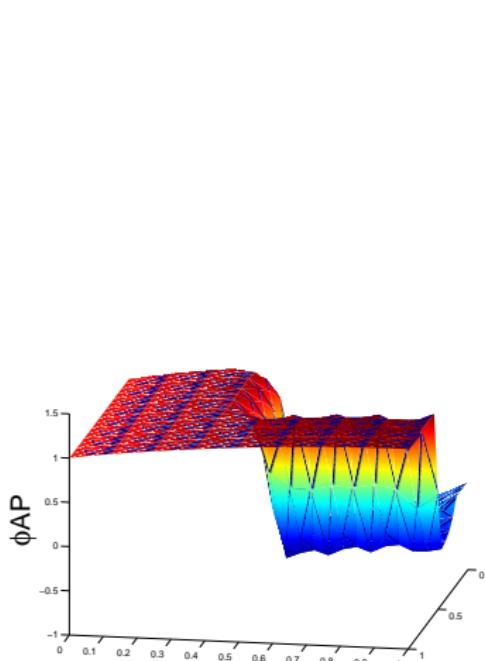
-relative errors in L^2 norms : On N : 0.0532,

On P : 0.0482,

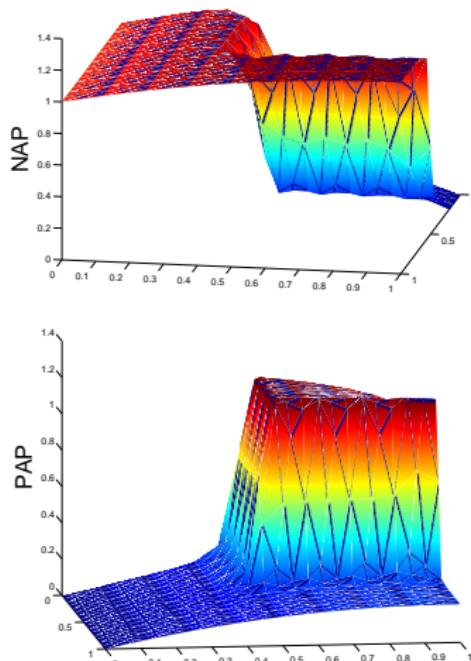
On ϕ : 0.0667.

PN Diode, $\lambda^2 = 10^{-10}$, $\Delta t = 5 \cdot 10^{-4}$

Implicit scheme is unstable



Our uncoupled AP scheme



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Works in progress

- ➡ Uniform convergence of implicit scheme (full discretization) and of our AP scheme.
- ➡ Study of the numerical initial and boundary layers.
- ➡ Long time behavior of the scheme
Semi-implicit scheme (C. Chainais-Hillairet, F. Filbet, 2010).
- ➡ Corrosion model of iron (C. Chainais-Hillairet, C. Bataillon)

$$\begin{cases} \varepsilon \partial_t N - \nabla \cdot (\nabla N - N \nabla \phi) = 0, \\ \partial_t P - \nabla \cdot (\nabla P + P \nabla \phi) = 0, \\ -\lambda^2 \Delta \phi = P - N + C. \end{cases}$$

Two small parameters ε and λ , build an AP scheme for the double limit.